A THEORETICAL MODEL OF THE THERMAL CONDUCTIVITY
OF IDEALIZED SOIL

by

Anil Misra, Ph.D.
Assistant Professor
Civil Engineering Department

Bryan R. Becker, Ph.D., P.E., Member ASHRAE
Associate Professor
Mechanical and Aerospace Engineering Department

Brian A. Fricke, Student Member ASHRAE
Research Assistant
Mechanical and Aerospace Engineering Department

University of Missouri-Kansas City
5605 Troost Avenue
Kansas City, MO  64110-2823
A Theoretical Model of the Thermal Conductivity of Idealized Soil

Anil Misra, Ph.D., Bryan R. Becker, Ph.D., P.E. (Member ASHRAE) and Brian A. Fricke (Student Member ASHRAE)
University of Missouri-Kansas City
5605 Troost Avenue, Kansas City, MO 64110-2823

ABSTRACT

Accurate prediction of soil thermal conductivity is of prime importance in the numerical simulation of heat transmission through soils. This paper focuses upon empirical and semi-empirical prediction methods for soil thermal conductivity. A family of empirical correlations are presented which relate soil thermal conductivity to saturation for five soil types: gravel, sand, silt, clay, and peat. These correlations are developed from a database of measured data available in the literature. Also, a theoretical model of soil thermal conductivity is developed for granular materials composed of rotund particles in an almost dry state. This theoretical model includes the effects of the micro-structure and the conductivity of the solid phase. It explicitly relates soil thermal conductivity to dry density and agrees well with experimental data. This paper also presents a review and discussion of those factors which affect soil thermal conductivity, previously reported prediction methods, and conductivity measurement techniques.

NOMENCLATURE

a Shape function used in Gemant’s equation
A cross-sectional area of a soil sample
b Shape function used in Gemant’s equation
B Function of dry density in Van Rooyen method
c Near-field pore parameter in Eq. (28)
C Function of mineral type in Van Rooyen method
d Far-field pore parameter in Eq. (28)
F Factor in De Vries correlation
G Shear modulus of a particle
g_a Shape factor in De Vries’ correlation
g_b Shape factor in De Vries’ correlation
g_c Shape factor in De Vries’ correlation
h Apex water
h_0 Water absorbed as a film around the soil particles
<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>H</td>
<td>Heat flux at a contact neighborhood</td>
</tr>
<tr>
<td>H(Ω)</td>
<td>Heat flux through the solid angle Ω</td>
</tr>
<tr>
<td>Hc</td>
<td>Heat flux through solid contact</td>
</tr>
<tr>
<td>Hm</td>
<td>Heat flux through the moisture bridge</td>
</tr>
<tr>
<td>Hm</td>
<td>Heat flux in the neighborhood of the m-th contact in Eq. (18)</td>
</tr>
<tr>
<td>k</td>
<td>Thermal conductivity of soil</td>
</tr>
<tr>
<td>ka</td>
<td>Thermal conductivity of air</td>
</tr>
<tr>
<td>kd</td>
<td>Soil thermal conductivity in the dry state</td>
</tr>
<tr>
<td>kf</td>
<td>Thermal conductivity of fluid phase</td>
</tr>
<tr>
<td>kij</td>
<td>Effective thermal conductivity tensor</td>
</tr>
<tr>
<td>ks</td>
<td>Thermal conductivity of solid phase (soil particles)</td>
</tr>
<tr>
<td>ks</td>
<td>Thermal conductivity of solid phase in the saturated state</td>
</tr>
<tr>
<td>kw</td>
<td>Thermal conductivity of water</td>
</tr>
<tr>
<td>Ke</td>
<td>Dimensionless function of soil saturation</td>
</tr>
<tr>
<td>n</td>
<td>Number density of particles</td>
</tr>
<tr>
<td>n</td>
<td>Normal vector at m-th contact in Eq. (18)</td>
</tr>
<tr>
<td>n</td>
<td>Normal vector at m-th contact in Eq. (18)</td>
</tr>
<tr>
<td>n(O)</td>
<td>Direction cosine of solid angle Ω</td>
</tr>
<tr>
<td>n(Ω)</td>
<td>Direction cosine of solid angle Ω</td>
</tr>
<tr>
<td>N</td>
<td>Total number of contacts in volume V</td>
</tr>
<tr>
<td>p</td>
<td>Coordination number (number of contacts per particle)</td>
</tr>
<tr>
<td>P</td>
<td>Porosity</td>
</tr>
<tr>
<td>q</td>
<td>Electric power input per unit length of probe</td>
</tr>
<tr>
<td>Q</td>
<td>Electric power input</td>
</tr>
<tr>
<td>r</td>
<td>Radius of contact area</td>
</tr>
<tr>
<td>R</td>
<td>Radius of particle</td>
</tr>
<tr>
<td>s</td>
<td>Function of granulometry in Van Rooyen method</td>
</tr>
<tr>
<td>S</td>
<td>Saturation</td>
</tr>
<tr>
<td>t1</td>
<td>Time</td>
</tr>
<tr>
<td>t2</td>
<td>Time</td>
</tr>
<tr>
<td>vs</td>
<td>Volume of a single particle</td>
</tr>
<tr>
<td>V</td>
<td>Volume of representative element in Eq. (18)</td>
</tr>
<tr>
<td>w</td>
<td>Moisture content</td>
</tr>
<tr>
<td>wu</td>
<td>Fractional volume of unfrozen water</td>
</tr>
<tr>
<td>xf</td>
<td>Volume fraction of fluid phase</td>
</tr>
<tr>
<td>xs</td>
<td>Volume fraction of solid phase</td>
</tr>
<tr>
<td>z</td>
<td>Shape function used in Gemant's equation</td>
</tr>
<tr>
<td>a</td>
<td>Ratio of mineral to air thermal conductivities</td>
</tr>
<tr>
<td>β</td>
<td>Heat flux parameter</td>
</tr>
<tr>
<td>δij</td>
<td>Kronecker delta</td>
</tr>
<tr>
<td>?T</td>
<td>Temperature difference or change</td>
</tr>
<tr>
<td>?X</td>
<td>Length of a soil sample</td>
</tr>
<tr>
<td>f</td>
<td>Coordinate angle</td>
</tr>
<tr>
<td>?</td>
<td>Non-dimensional parameter defined in Eq. (29)</td>
</tr>
<tr>
<td>?1</td>
<td>Coefficient in empirical soil thermal conductivity correlation</td>
</tr>
<tr>
<td>?2</td>
<td>Coefficient in empirical soil thermal conductivity correlation</td>
</tr>
<tr>
<td>?3</td>
<td>Coefficient in empirical soil thermal conductivity correlation</td>
</tr>
<tr>
<td>?4</td>
<td>Coefficient in empirical soil thermal conductivity correlation</td>
</tr>
<tr>
<td>?</td>
<td>Shape function used in Gemant's equation</td>
</tr>
<tr>
<td>?</td>
<td>Coordinate angle</td>
</tr>
</tbody>
</table>
I. INTRODUCTION

A. Motivation

Numerical simulation of heat transmission through soil requires accurate estimates of soil thermal conductivity. Soil thermal conductivity estimation methods that could be easily incorporated into computer models of heat transfer will have a wide variety of applications. The resultant prediction methods will be valuable in the calculation of heat loss through basements, slabs and crawl spaces, and will thereby improve the accuracy of hourly building energy analysis programs. Claridge (1988) reports that the impact of earth-contact heat transfer is comparable with that of infiltration. However, the lack of accessible and accurate foundation thermal analysis has prohibited realistic comparison of above and below grade energy losses, thus leading to sub-optimal energy conservation decisions.

Soil thermal conductivity estimation methods will also facilitate the efficient design of ground source heat pump systems by providing more accurate estimates of thermal energy transfer between the earth loop heat exchanger and the surrounding soil (Hart and Whidden 1984). This will lead to better selection of backfill materials, and optimization of backfill moisture content.

In addition, these estimation methods will assist in assessing the effects of berms on the heat loss from thermal energy storage tanks (Rosen and Hooper 1989). They will also be useful in the simulation of the operation of seasonal heat storage in various soil formations which are, by far, the least expensive and most widely available storage media (Doughty et al. 1983).

Furthermore, the resultant prediction methods will help ensure the safe design and location of underground storage facilities for nuclear and other hazardous waste materials. Likewise, they will be useful to designers of buried cables for electric power transmission.

Several empirical and semi-empirical prediction methods for soil thermal conductivity have been
proposed. However, most of these methods are not suitable because they apply to only certain soil types under specific temperature and moisture conditions. Although a large quantity of measured data is available for various soil types and conditions, very little effort has been made to establish a generally applicable, semi-empirical soil thermal conductivity prediction method.

B. Scope of this Paper

In this paper, the authors focus upon prediction methods for soil thermal conductivity. In Section II there is a description of the factors which affect soil thermal conductivity, while in Section III a discussion of existing prediction techniques is presented.

As a basis for the development of more generally applicable prediction methods, a database of measured soil thermal conductivity data was created from the existing literature. A description of the measurement techniques as well as the soil samples used to obtain these data is given in Section IV. Section V describes an empirical model of soil thermal conductivity, while Section VI presents a theoretical model.

II. FACTORS AFFECTING SOIL THERMAL CONDUCTIVITY

Several factors, such as moisture content, dry density, mineral composition, temperature, and texture, influence soil thermal conductivity (Kersten 1949, Penner et al. 1975, Salomone et al. 1984, Salomone and Kovacs 1984, Salomone and Marlowe 1989, Brandon and Mitchell 1989, Mitchell 1991). Moisture content has a most notable effect upon soil thermal conductivity. As moisture is added to a soil, a thin water film develops around the soil particles, which bridges the gaps in the soil. This "bridging" increases the effective contact area between the soil particles, which increases the heat flow and results in higher thermal conductivity. As more moisture is added, the voids between the soil particles become completely filled with moisture and the soil thermal conductivity no longer increases with increasing moisture content (Salomone et al. 1984, Salomone and Kovacs 1984, Salomone and Marlowe 1989). The point at which the maximum thermal conductivity occurs is called the critical moisture content. For most soils, the critical moisture content coincides with the optimum moisture content, which is also the state of maximum dry density of the soil. However, for low density clay soils, the critical moisture content coincides with the plastic limit.
Soil thermal conductivity also increases with the dry density of the soil. With an increase in the soil's dry density, more soil particles are packed into a unit volume and, thus, the number of contact points between the particles increases. This increase in contact points provides a larger heat flow path resulting in higher soil thermal conductivity (Salomone et al. 1984, Salomone and Kovacs 1984, Salomone and Marlowe 1989).

The mineral composition of a soil influences its thermal conductivity. For example, sands with a high quartz content generally have a greater thermal conductivity than sands with high contents of plagioclase feldspar and pyroxene (Kersten 1949). Also, soils with high organic content generally have lower soil thermal conductivity (Salomone and Marlowe 1989).

A dramatic change in soil thermal conductivity occurs at the freezing point when the primary mode of heat transfer changes from convection to conduction (Kersten 1949, Penner et al. 1975, Salomone and Marlowe 1989). However, in other temperature ranges the variation of soil thermal conductivity with temperature is minimal (Brandon and Mitchell 1989). For example, Kersten (1949) reports an average variation of 4% in soil thermal conductivity over the temperature range from 4.4°C to 21°C. At low moisture content, the thermal conductivity is higher in the unfrozen state than in the frozen state. However, the converse is true at high moisture content (Penner et al. 1975).

Soil texture is another factor which may influence a soil's thermal conductivity. For a given moisture content and dry density, the thermal conductivity of coarse textured, angular grained soils is higher than that of fine textured soils (Kersten 1949, Salomone and Marlowe 1989). Also, uniformly graded soils exhibit lower thermal conductivity than well graded soils (Salomone and Marlowe 1989).

III. EXISTING PREDICTION METHODS

Many correlations for estimating soil thermal conductivity have been proposed. Van Rooyen and Winterkorn (1957), Johansen (1975), De Vries (1952), Gemant (1952), and Kersten (1949), among others, have all developed correlations for soil thermal conductivity. These correlations vary in complexity, and
each method is limited to only certain soil types under specific conditions. Farouki (1986) has published a
survey of these various correlations which is summarized below.

Van Rooyen's correlation (Van Rooyen and Winterkorn 1957), based on data collected from different
types of sand and gravel, is given as follows:

\[
\frac{1}{k} = B 10^{CS} + s
\]  

In this expression, \(k\) is soil thermal conductivity, W/m-K; \(S\) is the degree of saturation; and, \(B\), \(C\), and \(s\) are
functions of dry density, mineral type, and granulometry, respectively. The Van Rooyen method is limited
to unfrozen sand and gravel with saturation levels between 1.5% to 10%.

Johansen's correlation (Johansen 1975), which is based on thermal conductivity data for dry and
saturated states at the same dry density, has the following form:

\[
k = (k_{sat} - k_d) \cdot K_e + k_d
\]

In this correlation, \(k\) is soil thermal conductivity, W/m-K; \(k_{sat}\) and \(k_d\) are soil thermal conductivity in the
saturated and dry states, respectively; and \(K_e\) is a dimensionless function of soil saturation.

Based on data from Kersten (1949), Johansen derived equations for \(K_e\) in terms of saturation, \(S\). For
coarse, unfrozen soil:

\[
K_e = 0.7 \log S + 1.0
\]

For fine, unfrozen soil:

\[
K_e = \log S + 1.0
\]

For any type of frozen soil, \(K_e = S\).

Johansen believed that in a dry state, soil thermal conductivity was based upon dry density or porosity.
For dry natural soils, he derived an equation for \(k_d\) in terms of dry density, \(\rho_d\), specifically:

\[
k_d = \frac{0.135}{2700} \rho_d + \frac{64.7}{0.947} \rho_d \pm 20\%
\]

For crushed rock material, Johansen found that:
which is a function of porosity, P.

In the saturated state, Johansen proposed a simple geometric mean to calculate the soil thermal conductivity. For unfrozen soils:

\[ k_{sat} = k_s^{(1-P)} \cdot k_w^P \]  

(7)

where \( k_w \) is the thermal conductivity of water, and \( k_s \) is the thermal conductivity of the soil particles. For saturated, frozen soils containing some unfrozen water:

\[ k_{sat} = k_s^{(1-P)} \cdot 2.2^P \cdot 0.269^w_u \]  

(8)

where \( w_u \) is the fractional volume of unfrozen water. Johansen's method is suitable for calculating soil thermal conductivity of both coarse and fine-grained soils in the frozen and unfrozen states. However, it is limited to saturations greater than 20%.

The correlation given by De Vries (1952) assumes that soil is a two-phase material composed of uniform ellipsoidal particles dispersed in a fluid phase. The De Vries correlation is given as:

\[ k = x_f \cdot k_f + \frac{F \cdot x_s \cdot k_s}{x_f + F \cdot x_s} \]  

(9)

In this expression, subscripts f and s denote the fluid and solid phases, respectively; \( x \) is the volume fraction; and \( k \) is soil thermal conductivity, W/m·°C. The factor \( F \) is given by:

\[ F = \frac{1}{3} \sum_i \left[ 1 + \left( \frac{k_{s_i}}{k_f} - 1 \right)^{g_i} \right], \quad i = a, b, c \]  

(10)

In Eq. (10), the \( g \) values, which sum to unity, were originally intended to be shape factors, but are usually used to fit empirical data. De Vries found that by assuming \( g_a \) and \( g_b \) to be 0.125, Eq. (9) would be within 10% of the experimental data for saturated soils having low \( k_s/k_f \) ratios. For dry soils, \( k_s/k_f \) ratio of approximately 100, Eq. (9) gives values about 25% too low. De Vries' method is applicable to unfrozen coarse soils with saturations between 10% to 20%.
Gemant's correlation (Gemant 1952) is based upon the idealized soil particle shown in Figure 1. Gemant assumed that the idealized particles made contact only at their apexes and that water collected around these contact points to form a thermal bridge. Also, heat flow was assumed to be vertically upward.

Gemant's correlation is given as follows:

\[
\frac{1}{k} = \frac{\left[ (1 - a)/a \right]^{1/3} \tan^{-1} \left[ (k_s - k_w)/k_w \right]^{1/2}}{(h/2)^{1/3} \left[ k_w (k_s - k_w) \right]^{1/2}} + \frac{(1 - z)}{k_s a} \psi \left( \frac{b^2}{a} \right)
\]

(11a)

\[
a = 0.078 \rho_d^{1/2}; \quad h = 0.16 \times 10^3 \rho_d w - h_0
\]

(11b,c)

\[
z = \left( \frac{1 - a}{a} \right)^{2/3} \left( \frac{h}{2} \right)^{1/3}; \quad b^2 = \left( \frac{a}{1 - a} \right)^{2/3} \left( \frac{h}{2} \right)^{2/3}
\]

(11d,e)

In Eq. (11), \( \rho_d \) is dry density of soil; \( w \) is moisture content; \( h \) is the apex water (water collected around the contact points); \( h_0 \) is water absorbed as a film around the soil particles; \( k_s \) is the thermal conductivity of the solids; \( k_w \) is the thermal conductivity of water; and \( a, b, z, \) and \( \psi(b^2/a) \) are shape functions given by Gemant (1952). Gemant's method gives reasonable results for unfrozen sandy soils only.

Kersten (1949) tested many soil types and based his correlations on the empirical data he collected. He produced equations for frozen and unfrozen silt-clay soils and sandy soils. Kersten's correlations for unfrozen and frozen silt-clay soils are as follows:

Unfrozen: \( k \) is soil thermal conductivity (W/m·°C); \( w \) is moisture content; and \( \rho_d \) is dry density (kg/m³).

\[
\text{Unfrozen} : \quad k = [0.130 \log w - 0.0288] 10^{0.000624 \rho_d}
\]

(12a)

\[
\text{Frozen} : \quad k = 0.0110 (10)^{0.000812 \rho_d} + 0.00462 (10)^{0.000911 \rho_d} w
\]

(12b)

The correlations for sandy soils are as follows:

Unfrozen: \( k \) is soil thermal conductivity (W/m·°C); \( w \) is moisture content; and \( \rho_d \) is dry density (kg/m³).

\[
\text{Unfrozen} : \quad k = [0.1011 \log w + 0.0577] 10^{0.000624 \rho_d}
\]

(13a)

\[
\text{Frozen} : \quad k = 0.0110 (10)^{0.000812 \rho_d} + 0.00462 (10)^{0.000911 \rho_d} w
\]

(13b)

In Eq. (12) and Eq. (13), \( k \) is soil thermal conductivity (W/m·°C); \( w \) is moisture content; and \( \rho_d \) is dry density (kg/m³). The equations for the silt and clay soils apply for moisture contents of 7% or more; those for the sandy soils, of 1% or more. Kersten's correlations give reasonable results only for frozen soils with
saturations up to 90%.

As summarized in Table 1, Farouki (1986) has studied the applicability of these methods and has suggested the conditions under which each method should be used. It is clear that these methods are applicable only for limited soil types and conditions, and they do not offer a unified, cogent technique for the estimation of soil thermal conductivity. General purpose numerical models of heat transfer through soil require unified prediction methods for soil thermal conductivity which are applicable to a wide variety of soil types. The intent of the study reported in this paper was to develop a prediction method for a variety of soil types and conditions.

IV. DATABASE OF THERMAL CONDUCTIVITY MEASUREMENTS

A. Measurement Techniques

Two testing techniques have been employed to determine the thermal conductivity of soils: the steady state method and the transient method. In steady state testing, a temperature gradient is imposed across a soil sample. When the temperatures within the sample stabilize, the power required to maintain the temperature gradient is used to determine the thermal conductivity by using the equation:

\[ k = \frac{Q}{A} \frac{\Delta X}{\Delta T} \]  

(14)

where Q is the electric power input, A is the cross-sectional area of the sample, \( \Delta X \) is the length of the sample and \( \Delta T \) is the temperature difference imposed on the sample. A drawback of this method is that the time required to achieve steady state may result in drying at the hot end of the sample, and, accumulation of water or ice at the cold end. This moisture migration may lead to inaccurate thermal conductivity measurement (Kersten 1949, Penner et al. 1975, Farouki 1986).

The transient testing technique involves inserting a thin heat probe into the soil sample. A constant electric power is applied to the heat probe and the variation in probe temperature over time is noted. With this information, the thermal conductivity of the soil sample can be estimated by using the equation:
where \( q \) is the electrical power input per unit length of the probe and \( \Delta T \) is the temperature change between the times \( t_2 \) and \( t_1 \). Due to the shorter testing time required, the transient testing technique is less prone to moisture migration. Thus, the thermal conductivity values measured with the transient test are more accurate than those from the steady state test provided that there is good thermal contact between the soil and the probe (Penner et al. 1975, Salomone et al. 1984, Salomone and Kovacs 1984, Salomone and Marlowe 1989, Farouki 1986).

B. Description of Database

In this study a database was created from measured data available in the literature. Thermal conductivities at various dry densities, moisture contents and temperatures were collected for various soil types. To obtain reasonable results, many sources of data were consulted: Kersten (1949), Penner et al. (1975), Salomone and Marlowe (1989), De Vries (1952), Farouki (1986), Andersland and Anderson (1978), Nakshabandi and Kohnke (1965), and Sawada (1975). The measured soil thermal conductivity data reported in these sources were obtained by performing either a steady state or a transient test. Of the data sources cited in this paper, only Kersten made use of the steady state test. Based upon texture, in this study the soil data were classified into five general types, namely, gravel, sand, silt, clay and peat. A brief description of the soil samples which constitute the database is given below for each of the five soil types.

B.1. Gravel

Most of the measured data on gravels is from Kersten (1949). This data includes Chena river gravel which is mainly composed of quartz and igneous rock with sizes ranging from 2.5 to 19.0 mm. B.2. Sand

The measured data on sand was collected from the works of Kersten (1949), Salomone and Marlowe (1989), De Vries (1952), Andersland and Anderson (1978), Nakshabandi and Kohnke (1965), and Sawada (1975).

Kersten (1949) presented data on 12 sand samples, of which five were natural sands and seven were man-made. The five natural sands include Fairbanks Sand, Lowell Sand, Northway Sand, Northway Fine Sand, and Northwest Sand. The seven man-made sands include Fairbanks Sand, Lowell Sand, Northway Sand, Northway Fine Sand, and Northwest Sand.
Sand, and Dakota Sandy Loam. The Fairbanks sand was a siliceous sand with 27.5% of the particles larger than 2.0 mm and 70% of the particles between 0.5 and 2.0 mm. The Lowell sand was also siliceous with particles between 0.5 and 2.0 mm. The two Northway Sands were similar in their composition, with their main constituent being feldspar, and having grain sizes ranging from 4.75 mm to 0.075 mm. No details are available on the Dakota Sandy Loam.

Of the seven man-made sands, three were feldspar sands and four were quartz sands. The feldspar sands consisted of 90% sand-sized particles and 10% gravel-sized particles. The quartz sands included one sample with grain size larger than 0.5 mm and three samples with grain sizes between 0.5 mm to 2.0 mm.

The sands tested by Salomone and Marlowe (1989) were classified according to the Unified Soil Classification System (USCS). These sands included well graded sands (SW), poorly graded sands (SP), silty sands (SM), and clayey sands (SC). However, no information was available concerning their mineral constituents.

The remaining sands were fine grained sands, however, no information is available on their grain size distributions or mineral constituents.

B.3. Silt

The measured data on silt is from Kersten (1949) and Salomone and Marlowe (1989). Kersten tested three silts: Northway Silt Loam, Fairbanks Silt Loam, and Fairbanks Silty Clay Loam. All three silts were classified as low plasticity silts (ML) according to the USCS. Salomone and Marlowe presented data for several low plasticity silts. Little information is available about the mineral constituents of these silts.

B.4. Clay

The measured data on clay is from Kersten (1949), Salomone and Marlowe (1989), and Penner et al. (1975). Kersten tested two clays, Ramsey Sandy Loam and Healy Clay, both of which were classified as low plasticity clays (CL). The main mineral constituent of these clays is kaolinite. Although Salomone and Marlowe tested both high and low plasticity clays, no information was given concerning the mineral composition of these clays. The clay samples tested by Penner et al. were low plasticity clays containing quartz, illite, chlorite and kaolinite.
B.5. Peat

The measured data on peat is from Kersten (1949) and Salomone and Marlowe (1989). Kersten tested Fairbanks Peat while Salomone and Marlowe tested highly decomposed woody peat.

V. AN EMPIRICAL MODEL OF THERMAL CONDUCTIVITY

Moisture content or saturation has a significant effect upon soil thermal conductivity. Herein we briefly describe correlations of soil thermal conductivity as a function of saturation which were developed by Becker et al. (1992).

As depicted in Figure 2, the thermal conductivity of a soil increases in three stages as the saturation level increases. At low saturations, moisture first coats the soil particles. The gaps between the soil particles are not filled rapidly and thus there is a slow increase in thermal conductivity. When the particles are fully coated with moisture, a further increase in the moisture content fills the voids between particles. This increases the heat flow between particles, resulting in a rapid increase in thermal conductivity. Finally, when all the voids are filled, further increasing the moisture content no longer increases the heat flow and the thermal conductivity does not appreciably increase. An empirical model used to describe this behavior is as follows:

\[
S = \lambda_1 \left[ \sinh (\lambda_2 k + \lambda_3) - \sinh (\lambda_4) \right] \tag{16}
\]

In Eq. (16), \(S\) is the saturation; \(k\) is the soil thermal conductivity (W/m·°C); and \(\lambda_1, \lambda_2, \lambda_3, \) and \(\lambda_4\) are coefficients which are specified in Table 2 for each of the five soil types in both the frozen and unfrozen states. At a saturation of zero, Eq. (16) reduces to the following:

\[
\lambda_2 k_d + \lambda_3 = \lambda_4 \tag{17}
\]

Eq. (17) shows that the coefficient \(\lambda_4\) is related to the thermal conductivity of dry soil, \(k_d\).

Figures 3 through 7 present the measured soil thermal conductivity versus saturation data for the five soil types in both the frozen and unfrozen states. The empirical correlations, based upon Eq. (16), are also plotted on Figures 3 through 7. Three curves have been given for each soil type (except peat). The middle
curve represents the mean of the measured data. The upper and lower curves approximate the range of measured data. Due to the small amount of measured data for peaty soils, only a mean correlation is presented in Figure 7. Measured data collected for gravel includes saturations up to approximately 40% and thus, the correlations for gravel are valid only to 40% saturation. An error analysis of these correlations is presented in the work of Becker et al. (1992).

VI. A THEORETICAL MODEL OF THERMAL CONDUCTIVITY

Very few attempts have been made to develop mathematical models of soil thermal conductivity which account for the micro-structure of the soil. Such development has been impeded by the complex nature of the heat transfer mechanism in a particulate system which is composed of two or three phases. However, an analysis can be made of a simpler particulate system, composed of random arrays of identical spheres in an almost dry state, as depicted in Stages 1 and 2 of Figure 2. The heat transfer in this simple system can be idealized to occur in the neighborhood of the inter-particle contacts. Based upon this idealized model, Batchelor and O'Brien (1972), derived the following expression for the effective thermal conductivity tensor, \(k_{ij}\), of a representative element of random arrays of identical spheres:

\[
k_{ij} = \frac{3 v_s}{2 V} k_a \sum_{m} H^{m}_{i,j} \quad (i, j = 1, 2, 3; m = 1, 2, 3, \ldots N)
\]

Here, \(v_s\) is the volume of a single particle; \(V\) is the total volume of the representative element; \(k_a\) is the thermal conductivity of air; \(H^{m}_{i,j}\) is the non-dimensional heat flux in the neighborhood of the \(m\)-th contact; \(n\) and \(n\) are the normal vectors at the \(m\)-th contact for the \(n\)-th sphere; and \(N\) is the total number of contacts in the representative element. Throughout this section, tensor convention is assumed for all subscripts unless noted otherwise.

Eq. (18) applies to a granular packing with a finite number of contacts. For a representative volume with a large number of contacts, it is useful to express the effective thermal conductivity tensor in a continuum form. Introducing a distribution density function of inter-particle contacts as \(\Omega(\Theta)\), the number of contacts in the solid angle interval \(\Theta\) to \(\Theta + d\Theta\) is given by \(N \Omega(\Theta)d\Theta\), where \(N\) is the total number of
contacts in the representative element of volume $V$. Thus, the effective thermal conductivity tensor, given in Eq. (18), can be written in an integral form as follows:

$$k_{ij} = \frac{3 v_s N}{2V} k_s \int_{\Omega} n_i(\Omega) n_j(\Omega) H(\Omega) \xi(\Omega) d\Omega$$  \hspace{1cm} (19)$$

where $n_i(\Omega)$ and $n_j(\Omega)$ are the direction cosines of the solid angle $\Omega$, and $H(\Omega)$ is the non-dimensional heat flux through the solid angle $\Omega$. In Eq. (19), the integration is performed as follows:

$$\int_{\Omega} (\cdot) d\Omega = \int_0^{2\pi} \int_0^\pi (\cdot) \sin \phi d\phi d\theta$$  \hspace{1cm} (20)$$

where $d\Omega$ is the solid angle, and $\phi$ and $\psi$ are coordinate angles, as shown in Figure 8. Assuming that the heat flux, $H$, at all contact neighborhoods is equal, Eq. (19) can be simplified to:

$$k_{ij} = \frac{3 v_s N}{2V} k_s H \int_{\Omega} n_i(\Omega) n_j(\Omega) \xi(\Omega) d\Omega$$  \hspace{1cm} (21)$$

For isotropic distribution of contacts, the distribution function, $\xi(\Omega)$, is given as follows (Misra 1992):

$$\xi(\Omega) = \frac{1}{4\pi}$$  \hspace{1cm} (22)$$

In this case, the effective conductivity tensor becomes:

$$k_{ij} = \frac{v_s N}{2V} k_s H \delta_{ij}$$  \hspace{1cm} (23)$$

Noting that $v_s N/V = p_n$, Eq. (23) becomes:

$$k = \frac{1}{2} v_s p_n k_s H$$  \hspace{1cm} (24)$$

where $\delta_{ij}$ is the Kronecker delta ($\delta_{ij}=1$ for $i=j$; $\delta_{ij}=0$ for $i \neq j$), $p$ is the coordination number (number of contacts per particle) and $n$ is the number of particles per unit volume. In Eq. (24) the terms $v_s$, $p$ and $n$ are measures of the packing density of a granular assembly. It is useful to replace these terms by more direct measures of granular media density, namely, the dry density, $\rho_d$, and the solid density, $\rho_s$. An expression for the coordination number, $p$, can be obtained from the void ratio-coordination number relationship reported
by Chang et al. (1990) and the product $v_n$ can be replaced by $\frac{d}{s}$ to yield the following:

$$p = \left( 21 - 8 \frac{\rho_s}{\rho_d} \right)$$  \hspace{1cm} (25)

Thus Eq. (24) becomes:

$$k = \left( 11 \frac{\rho_d}{\rho_s} - 4 \right) k_s H$$  \hspace{1cm} (26)

In Eq. (26), the key requirement is the heat flux, $H$, at a contact neighborhood. Batchelor and O'Brien (1972) studied the heat transfer through two particles in contact. Based upon their study, an approximate expression for the non-dimensional heat flux, $H$, is written as follows:

$$H = H_c(\lambda) + w H_m(\lambda) + 2 \ln \alpha + \beta$$  \hspace{1cm} (27)

where $H_c$ is the flux through the solid contact; $H_m$ is the flux through the moisture bridge in the neighborhood of contact; $w$ is the moisture content; and $\alpha$ is the ratio of mineral to air thermal conductivities, i.e. $k_m/k_a$. The non-dimensional heat flux parameter $\beta$ may be written as follows:

$$\beta = d + w c - 3.9$$  \hspace{1cm} (28)

Here, $d$ depends upon conditions within the pores far from the contact; while $c$ depends upon conditions near the contact and is more sensitive to the presence of moisture. In Eq. (27), $\lambda$ is a non-dimensional parameter, which is defined as follows:

$$\lambda = \frac{\alpha r}{R}$$  \hspace{1cm} (29)

Here, $R$ is the radius of the particle and $r$ is the radius of the contact area. The contact radius depends upon the loading conditions on the granular medium. For example, under isotropic loading conditions, the contact radius may be obtained from Hertz-Mindlin contact theory as follows (Chang and Misra 1990):

$$r = \left[ \frac{3\pi \rho_s (1 - \nu)}{2p \rho_d G} \sigma_o \right]^{1/3} R$$  \hspace{1cm} (30)
Here, $G$ is the shear modulus of a particle; $\nu$ is the Poisson's ratio of the particle; and $s_o$ is the isotropic confining stress. Batchelor and O’Brien (1972) have reported simple approximations for fluxes $H_c$ and $H_m$ depending upon the value of $\lambda$, given as follows:

$$H_c = 0.22 \lambda^2; \quad H_m = -0.05 \lambda^2$$  \hspace{1cm} (31)

for $\lambda < 1$, and

$$H_c = \frac{2 \lambda}{\pi}; \quad H_m = -2 \ln \lambda$$  \hspace{1cm} (32)

for $\lambda > 1$.

By combining Eq. (27) through Eq. (31), it can be seen that as the confining stress, $s_o$, approaches zero, the heat flux contributions, $H_c$ and $H_m$, become negligible, and the expression for thermal conductivity may be simplified to:

$$k = k_a \left( 11 \frac{\rho_d}{\rho_s} - 4 \right) \left( 2 \ln \alpha + d + wc - 3.9 \right)$$  \hspace{1cm} (33)

Eq. (33) shows that thermal conductivity varies linearly with dry density. The measured data exhibits a similar behavior. Also, Eq. (33) accounts for the presence of moisture at low levels. This behavior is similar to the experimental results reported by Brandon and Mitchell (1989) for sands with low moisture content. In Figures 9 and 10, a comparison is presented between the predictions and the measurements for thermal conductivity versus dry density for quartz and feldspar sands at three levels of moisture content: $w = 0.0\%$, 0.5\%, and 1.0\%. In these predictions, the thermal conductivity of a solid particle is taken to be 8.4 W/m·°C and that of air is taken to be 0.026 W/m·°C. The parameter, $d$, is set to -3.1 for quartz and -3.9 for feldspar, and $c$ is set to 10.5 for quartz and 6.5 for feldspar. As shown in these figures, the model successfully predicts the same trends as that exhibited by the measured data discussed in Section IV.

VI. CONCLUSIONS

This paper has focused upon prediction methods for soil thermal conductivity. A family of empirical
correlations has been presented which relates soil thermal conductivity to saturation for five soil types, namely, gravel, sand, silt, clay, and peat, in both the frozen and unfrozen states. These empirical prediction methods were developed from a database of soil thermal conductivity measurements available in the literature. The unified format of these correlations make them readily adaptable to numerical heat transfer algorithms.

Also, a theoretical model of soil thermal conductivity has been presented which is based upon a micro-structural approach. Although this model considers a rather simplified particulate system, it gives insight into the mechanism by which dry density influences soil thermal conductivity. At present, this model is limited to granular materials in an almost dry state. In spite of this limitation, the model correctly exhibits the effects of the presence of moisture at low levels.

This paper has also presented a review and discussion of factors which affect soil thermal conductivity, previously reported prediction methods, and conductivity measurement techniques.
REFERENCES


Table 1.

Applicability of Prediction Methods

<table>
<thead>
<tr>
<th>State</th>
<th>Texture</th>
<th>Saturation</th>
<th>Method</th>
</tr>
</thead>
<tbody>
<tr>
<td>Unfrozen</td>
<td>Coarse Grained</td>
<td>0.015 - 0.100</td>
<td>Van Rooyen (except for low-quartz crushed rock)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.100 - 0.200</td>
<td>De Vries</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.200 - 1.000</td>
<td>Johansen</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.000 - 1.000</td>
<td>Gemant (sandy silt-clay)</td>
</tr>
<tr>
<td>Fine Grained</td>
<td></td>
<td>saturated</td>
<td>Johansen, De Vries, Gemant</td>
</tr>
<tr>
<td>Frozen</td>
<td>Coarse Grained</td>
<td>0.100 - 1.000</td>
<td>Johansen</td>
</tr>
<tr>
<td></td>
<td></td>
<td>saturated</td>
<td>Johansen, De Vries</td>
</tr>
<tr>
<td>Fine Grained</td>
<td></td>
<td>0.000 - 0.900</td>
<td>Kersten</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.100 - 1.000</td>
<td>Johansen</td>
</tr>
<tr>
<td></td>
<td></td>
<td>saturated</td>
<td>Johansen, De Vries</td>
</tr>
</tbody>
</table>

aData from Farouki (1986)
TABLE 2

Correlation Coefficients

<table>
<thead>
<tr>
<th>Soil Type</th>
<th>Frozen</th>
<th>Unfrozen</th>
<th>( ?_1 )</th>
<th>( ?_2 )</th>
<th>( ?_3 )</th>
<th>( ?_4 )</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Low</td>
<td>Mean</td>
<td>High</td>
<td>Low</td>
<td>Mean</td>
<td>High</td>
</tr>
<tr>
<td>Clay</td>
<td>Frozen</td>
<td>23.5</td>
<td>14.5</td>
<td>14.0</td>
<td>1.73</td>
<td>1.73</td>
</tr>
<tr>
<td></td>
<td>Unfrozen</td>
<td>33.5</td>
<td>27.0</td>
<td>14.0</td>
<td>2.01</td>
<td>1.84</td>
</tr>
<tr>
<td>Gravel</td>
<td>Frozen</td>
<td>25.4</td>
<td>11.0</td>
<td>11.3</td>
<td>2.01</td>
<td>2.43</td>
</tr>
<tr>
<td></td>
<td>Unfrozen</td>
<td>16.5</td>
<td>6.5</td>
<td>8.3</td>
<td>2.22</td>
<td>2.63</td>
</tr>
<tr>
<td>Peat</td>
<td>Frozen</td>
<td>12.0</td>
<td></td>
<td></td>
<td>2.77</td>
<td></td>
</tr>
<tr>
<td></td>
<td>Unfrozen</td>
<td>28.0</td>
<td></td>
<td></td>
<td>6.00</td>
<td></td>
</tr>
<tr>
<td>Sand</td>
<td>Frozen</td>
<td>26.0</td>
<td>10.0</td>
<td>15.0</td>
<td>1.84</td>
<td>1.66</td>
</tr>
<tr>
<td></td>
<td>Unfrozen</td>
<td>6.4</td>
<td>6.8</td>
<td>6.8</td>
<td>5.55</td>
<td>2.77</td>
</tr>
<tr>
<td>Silt</td>
<td>Frozen</td>
<td>38.0</td>
<td>19.5</td>
<td>18.5</td>
<td>1.66</td>
<td>1.87</td>
</tr>
<tr>
<td></td>
<td>Unfrozen</td>
<td>28.0</td>
<td>17.0</td>
<td>22.0</td>
<td>2.77</td>
<td>2.77</td>
</tr>
</tbody>
</table>